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Highlights

- Analytical formulation of Wheatley surfaces containing two leaflets.
- Proposition of a non-linear model capable of capturing the mechanical behavior of the valve under simplified service conditions.
- Taking non-linear contact forces between leaflets into account for equilibrium stability.
- The results show that the leaflet height does not change the mechanism of force transfer along the valve.

Mathematical representation and nonlinear modelling of the Wheatley mitral valve

H. L. Oliveira^{a,b,*}, G. C. Buscaglia^b, J. A. Cuminato^b, S. McKee^c, I. W. Stewart^c, M. M. Kerr^d, D. J. Wheatley^e

^aFECFAU - Departamento de Estruturas, Universidade Estadual de Campinas, Cidade Universitária, Av. Albert Einstein, 901, 13083-852, SP, Brazil

^bInstituto de Ciências Matemáticas e de Computação - ICMC, Universidade de São Paulo, Campus de São Carlos, Caixa Postal 668, 13560-970, SP, Brazil

^cDepartment of Mathematics and Statistics, University of Strathclyde, 26 Richmond St, Glasgow, G1 1XH, United Kingdom

^dDepartment of Biomedical Engineering, University of Strathclyde, 50 George St, Glasgow, G1 1QE, United Kingdom

^eSchool of Science and Engineering, University of Dundee, Nethergate, Dundee, DD1 4HN, United Kingdom

Abstract

This study is concerned with the Wheatley design of the mitral valve. A mathematical description, in terms of elementary functions, is provided for the S-shaped leaflets. This is based on a level set containing symmetric circles (or more generally ellipses) which allow parametrization. A geometric nonlinear mechanical model subjected to a uniform pressure gradient and in the absence of inertial forces is introduced. The model results in a system of nonlinear equations that is solved using iterative incremental techniques. Under normal pressure loads, the S-shaped geometries induce internal forces which manifest themselves in two combined effects: bending and torsion. As a consequence, the supports are subject to periodic bending actions that tend

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^{*}Corresponding author

Email address: hugo.oliveira@unicamp.br (H. L. Oliveira)

to deform the support frame towards the interior of the valve. Providing resistance becomes vital for maintaining stable equilibrium. It is also observed that for circular base shape geometries, the mechanism for transmitting the equilibrium forces remains unchanged when the height/diameter ratio is kept below 2.

Keywords: Leaflets analytical description, Finite Element Method, nonlinear analysis, Static structural model

1 1. Introduction

The human heart has a set of four values that open and close to ensure 2 that blood flows in the correct direction. The four valves are the aortic valve, mitral valve, pulmonary valve and tricuspid valve. Some conditions such as ageing, congenital heart disease, infections, diabetes among others can lead 5 to malfunction of these values affecting the health of the individual [1]. There 6 are two main types of heart problems that can happen in any of the valves: 7 stenosis, when the valve does not open properly and regurgitation, when 8 the blood flows back due to poor valve closure. Diseases associated with 9 the aortic value account for 61% of deaths from heart value disease, while 10 diseases associated with the mitral valve account for 15%, with regurgitation 11 being one of the most prevalent occurrences [2]. 12

In general, valve pathologies can be divided into two classes: primary (or degenerative), when the disease is related to the valve itself, and secondary, when the disease results from dysfunction of adjacent cardiac structures. In the case of mitral regurgitation, primary causes are leaflet prolapse or chordal rupture, and secondary causes are ischaemic heart disease or cardiomyopathy

[3]. When the mitral valve functions poorly and there is no possibility of
¹⁹ surgical repair, then valve replacement becomes a treatment option [4].

There are two main types of replacement mitral valves: tissue valves 20 and mechanical valves. Tissue valves, also called biological or bioprosthetic 21 valves, are made from animal tissue. In general, tissue valves wear out with 22 time and require new medical interventions. On the other hand, mechani-23 cal valves are made from durable materials and generally do not wear out. 24 However, patients will need to take anticoagulants for life [5, 6]. As a con-25 sequence, the development of replacement valves that are both durable and 26 minimise the impact on the patient's life is of prime importance, and remains 27 an active field of research [7, 8]. 28

In terms of valve replacement design, numerous studies have demon-29 strated the benefits that numerical and computer simulations can provide, 30 e.g. [9]. The formulations allow complex behaviours such as durability, bi-31 ological response and haemodynamics to be investigated in a single model. 32 From an engineering point of view, an ideal valve should not only cause 33 reduced pressure drops, but also ensure minimal regurgitation volume, min-34 imise turbulence production, avoid zones of high shear stresses and flow stag-35 nation [10]. In addition, the designer should avoid valve shapes that lead 36 to low washout performance since this is a coagulation-facilitating factor 37 [11, 12, 13, 14].38

Substantial work has been done by Wheatley et al. [15] on the function of experimental polymeric valves of conventional design. Intending to create a design that not only preserves the good features of existing prosthetic heart valves, but also potentially improves the washout effect, D. J. Wheatley

introduced the concept of S-shaped leaflets [16, 17, 18]. The geometry of the 43 leaflets can be obtained from the union of circular and elliptical contour lines 44 so that a family of shapes becomes readily obtainable, which is of crucial 45 importance for manufacturing processes [19, 20]. From these geometries, 46 computational models can be built to simulate the global behaviour of the 47 assembly [21]. Approaches such as these have been increasingly applied in 48 the design of state-of-the-art valves [22, 23, 24]. In the present study we 49 investigate Wheatley's design for bi-leaflet mitral valves. 50

The contribution of this study is twofold. We first present a formulation 51 describing the leaflet surface inspired by Wheatley's design with a view to 52 mitral prosthetic valve applications. Although the defining parameters are 53 different, it is essentially a natural extension of previous work on the Wheat-54 ley aortic valve where we introduced mathematical representations of the 55 Wheatley valve using elementary functions [19, 20, 21]. Secondly, we present 56 a geometric nonlinear structural model capable of predicting the behaviour of 57 the leaflets in the absence of inertial forces under uniform pressure gradients. 58

Producing a mathematical representation of the design is a useful first 59 step before applying computational fluid dynamics to understand both the 60 blood flow in the valve and the cavities (atrium, ventricle), and the internal 61 leaflet stresses during valve function. Of course, the shape could have been 62 generated via splines, but the principal advantage of a mathematical repre-63 sentation with elementary functions is that it permits shape changes to be 64 easily trialled. For instance, the three-dimensional shape can be obtained 65 directly by making linear increments along the z direction. However, this is 66 not the only option; quadratic, cubic or higher order increments and even 67

transcendental functions are also possible in an analytical approach. That
creates a family of designs, providing a flexibility that is not available with
the use of splines.

71 2. Numerical Modelling

72 2.1. Wheatley geometric design

⁷³ In this section we present a closed mathematical representation of the⁷⁴ Wheatley Mitral Valve.

75 2.1.1. The artificial mitral valve

A paper version of the mitral value is displayed in Figure 1. From this, we are able to deduce the contour lines (set of points where a given scalar function assumes a constant value - see Figure 2). Indeed, close inspection of the paper value and its associated contour lines suggests that, for each contour line, there exists four underlying symmetrically placed circles. These are displayed in Figure 3; the great circle is, for convenience, and without any loss of generality, the unit circle.

⁸³ 2.1.2. Generalization from circles to ellipses

In this section we shall, nonetheless, replace the circles with ellipses: not only is this more general providing an infinite number of designs, it also allows us to obtain the original Wheatley design as a special case. Thus, in place of Figure 3, consider Figure 4 and the two ellipses, CHB, CFD lying on the axis y = 0. Let their major axes be 2a and 2(1 - a) (so that their foci are a and 1 - a), respectively. Furthermore, let them both pass through x = 1.



Figure 1: Paper model of the Wheatley mitral valve.

⁹¹ Their equations can readily be found to be :

$$\frac{(x - (1 - a))^2}{a^2} + \frac{y^2}{(\gamma a)^2} = 1,$$
(1)

$$\frac{(x-a)^2}{(1-a)^2} + \frac{y^2}{(\gamma(1-a))^2} = 1,$$
(2)

where γ is a measure of 'ellipticity' and is related to the common eccen-



Figure 2: Mitral valve contour lines.

 $_{93}$ tricity, e, of the two ellipses:

$$\gamma^2 = 1 - e^2.$$

To obtain the equations of the two ellipses AGD and AEB it is only necessary to rotate (1) and (2) by π radians. We therefore immediately obtain

$$\frac{(x+(1-a))^2}{a^2} + \frac{y^2}{(\gamma a)^2} = 1,$$
(3)

$$\frac{(x+a)^2}{(1-a)^2} + \frac{y^2}{(\gamma(1-a))^2} = 1.$$
(4)

The choice of $\gamma = 1$ reduces the above equations (1) ...(4) to those of \dots



Figure 3: Two large and two smaller circles with the auxiliary unit circle.

- ⁹⁸ circles shown in Figure 3.
- 99 2.1.3. Parametrization of the elliptic equations

For graphical purposes we require a parametrization of the four ellipses. Note we wish to trace the arcs CHBEA and AGDFC (red lines in Figure 4). We may write the four equations (1), (2), (3) and (4) in parametrized form as



Figure 4: One large circle and four ellipses.

$$x = 1 - a + a\cos(\theta), \ y = \gamma a\sin(\theta), \tag{5}$$

$$x = a + (1 - a)\cos(\theta), \ y = \gamma(1 - a)\sin(\theta), \tag{6}$$

$$x = -(1-a) + a\cos(\theta), \ y = \gamma a\sin(\theta), \tag{7}$$

$$x = -a + (1-a)\cos(\theta), \ y = \gamma(1-a)\sin(\theta).$$
(8)

Recall that we wish to trace out certain arcs: the arc CHB and the arc BEA, and the arc AGD and the arc DFC. In order to do this we require both starting and end points. These are the following: for (5) we need θ to range from 0 to π ; for (6) we require θ to go from π to 2π ; for (7) we need θ run from 0 to $-\pi$; and for (8) we must have θ ranging from 0 to π .

109 2.1.4. Examples of the Wheatley elliptic mitral valve

In this subsection we display examples of the geometrical representation of the generalized Wheatley mitral valve, for $\gamma = 1$ (Figure 5a) and $\gamma = \frac{1}{2}$ (Figure 5b). In the special case when $\gamma = 1$ (the original Wheatley valve) the ellipses reduce to circles. In this case the valve is displayed partially open (Figure 5c) and closed (Figure 5d).

The equations that describe the three-dimensional shape of the leaflets allow opening and closing to be performed. Although these motions are kinematically compatible, they do not preserve linear and angular momentum, nor do they follow material-specific responses. To be consistent with the first principles of mechanics, it is necessary to design a geometric nonlinear model, which will be detailed in the next subsection.

121 2.2. Computational mesh

In this step we establish a finite element mesh that describes the geometry 122 of interest. The original circular design ($\gamma = 1$) is chosen. The dimensions 123 assumed are: $20 \, mm$ diameter at the base, $10 \, mm$ high. At the top of the 124 leaflets there are two semi-circles of 5 mm radius. The mesh (Figure 6) 125 contains 19321 linear interpolation shell (SHELL181) elements (quadrilateral, 126 triangles). 200 μm element size and 19452 nodes. Each node has 6 degrees of 127 freedom (three translations and three rotations). Both leaflets have a uniform 128 thickness of $250 \,\mu m$. 129

130 2.3. Material response

In the present study, St Venant-Kirchhoff hyperelastic material is assumed. Its relevance lies not only in its simplicity, but also by serving as a



Figure 5: Mitral value for (a) $\gamma = 2$, (b) $\gamma = 1/2$ (c) $\gamma = 1$ (open) (d) $\gamma = 1$ (closed).

gateway to more elaborate material models. The resulting elasticity matrix
has the same form as the linear elasticity formulation, except that GreenLagrange strains are used. The second Piola-Kirchhoff stress tensor (S) is
symmetric and is related to the Green-Lagrange strain tensor (E) according



Figure 6: Finite element mesh in the reference configuration: (a) Top view and (b) Perspective view. (116.712 degrees of freedom)

137 to the St. Venant constitutive expression:

$$\mathbf{S} = \lambda \mathrm{tr}(\mathbf{E})\mathbf{I} + 2\mu \mathbf{E},\tag{9}$$

where λ and μ are material parameters known as Lamé constants, and I is the second order unit tensor. The quantities λ and μ are related to the usual Young Modulus (E) and Poisson ratio (ν) as follows:

$$\lambda = \frac{E\nu}{(1+\nu)(1-2\nu)},\tag{10}$$

141

$$\mu = \frac{E}{2(1+\nu)}.\tag{11}$$

The values employed in the model are shown in Table 1. It corresponds to a thermoplastic polyurethane (TPU), although the definitive Wheatley Valve material has not yet been established. The motivation for these parameters came from a previous study carried out by our group, which is described in
reference [24].

It is worth noting that this formulation takes into account both the potential initial stresses and the change of configuration due to large displacements.

Parameter	Description	Value
ν	Poisson's ratio	0.49
E	Young modulus	$65 \mathrm{MPa}$

Table 1: Input parameters for the mechanical models

149 2.4. Boundary conditions

When the valve is positioned between the atrial and ventricular cavities, the movement of the leaflets is induced by the pressure gradient exerted by the blood. In this study, this condition will be approximated assuming that the valve motion will be caused by a normal pressure evenly distributed along the surface following the variation shown in Figure 7. The base of the valve (which is annular) and the vertical posts are kept fixed throughout the stress analysis.

157 2.5. Contact conditions

¹⁵⁸ Contact modelling is crucial for predicting the functioning of the valve. ¹⁵⁹ In biomedical community, the contact region between the two fully closed ¹⁶⁰ leaflets is known as the coaptation surface. In pre-transplant patients, this ¹⁶¹ surface distributes the mechanical forces between the two leaflets and the ¹⁶² commissural regions of the valve. Proper coaptation modelling is important



Figure 7: Normal pressure load curve.

not only to obtain force distributions in agreement with first mechanical
principles but also to capture physiological haemodynamics.

As both surfaces are flexible, it makes no difference which leaflet will be 165 set as the *target*. We assume that the surfaces in contact can transmit shear 166 forces in addition to compressive forces. When the equivalent shear stress is 167 less than a limit frictional stress (τ_{lim}), no motion occurs between the two 168 surfaces. This state is known as sticking. Once the equivalent frictional stress 169 exceeds τ_{lim} , both surfaces will slide relative to each other. This condition is 170 known as sliding. When cohesion is assumed, sliding resistance exists even 171 in the absence of normal surface pressures. The transition point between 172 sticking and sliding is calculated according to Coulomb's Law: 173

$$\tau_{lim} = \alpha p + \beta, \tag{12}$$

$$||\tau|| = \sqrt{\tau_1^2 + \tau_2^2} \le \tau_{lim}.$$
 (13)

where τ_1 and τ_2 are frictional components in direction 1 and 2 (mutually perpendicular) defined on the contact surface, α is the friction coefficient, pis the contact normal pressure and β is the contact cohesion.

It is also possible to define the maximum equivalent frictional stress τ_{lim} so that, regardless of the magnitude of the contact pressure, sliding will occur if the magnitude of the equivalent frictional stress reaches this value [25]. Figure 8 displays a graphical representation of the frictional model.



Figure 8: Graphic representation of the friction model.

181 2.6. Computational Structural Mechanics

The general dynamical analysis of the mitral leaflets gives rise to the equation [26, 27]:

$$[\mathbb{M}]\{\ddot{U}\} + [\mathbb{C}]\{\dot{U}\} + [\mathbb{K}]\{U\} = \{F\},$$
(14)

where $\{U\}$ is the displacement, while $\{\ddot{U}\}$ and $\{\dot{U}\}$ represent the dis-184 placements' first and second derivatives, $[\mathbb{M}]$ is the mass matrix, $[\mathbb{C}]$ and $[\mathbb{K}]$ 185 are the material damping and stiffness matrices, respectively, and $\{F\}$ rep-186 resents the total forces, internal and external, acting on the mitral leaflets. 187 In this study, we assume that inertial and damping forces can be neglected. 188 The finite element mesh (Section 2.2) is used to build a displacement 189 field $(\{u\})$ approximated by polynomial shape function ($[\mathbb{H}]$ in matrix form), 190 $\{u\} = [\mathbb{H}]\{U\}$. Following Sung and Kwak [28], taking into account the 191 contact forces, equation (14) can be linearized to provide $\Delta\{U\}$, which reads 192 as follows: 193

$$([\mathbb{K}_L] + [\mathbb{K}_{NL}])\Delta\{U\} = \{F\}_{ext}^{t+\Delta t} - \{F\}_{int}^{t} + \{F\}_c^{t+\Delta t},$$
(15)

194 where

$$\{F\}_{ext}^{t+\Delta t} = \mathbf{A} \int_{\Omega_{0_e}} [\mathbb{H}]^\top \{b\} \ \mathrm{d}\Omega + \mathbf{A} \int_{\partial\Omega_{0_e}} [\mathbb{H}]^\top \{P\}_0 \ \mathrm{d}\partial\Omega, \tag{16}$$

$$\{F\}_{int}^{t} = \mathbf{A} \int_{\Omega_{0_{e}}} [\mathbb{B}_{L}]^{\top} [\mathbf{S}]^{t} \, \mathrm{d}\Omega, \qquad (17)$$

$$\{F\}_{c}^{\mathbf{t}+\Delta\mathbf{t}} = \mathbf{A} \int_{\Gamma_{0_{e}}^{c}} [\mathbb{H}]^{\top} (p\{\mathbf{n}\} + \tau\{\mathbf{m}\}) \mathrm{d}\Gamma_{0}^{c}, \tag{18}$$

$$[\mathbb{K}_{L}] = \mathbf{A} \int_{\Omega_{0_{e}}} [\mathbb{B}_{L}]^{\top} [\mathbb{D}] [\mathbb{B}_{L}] \, \mathrm{d}\Omega,$$

$$[\mathbb{K}_{NL}] = \mathbf{A} \int_{\Omega_{0_{e}}} [\mathbb{B}_{NL}]^{\top} [\mathbb{S}] [\mathbb{B}_{NL}] \, \mathrm{d}\Omega.$$
(20)

In these expressions, the superscript ' \top ' denotes the transpose operation, 195 't' and 't + Δt ' represents the previous and current time step and subindex 196 'e' indicates a given element. A denotes the standard assembling operator 197 over all elements in the mesh, $\{b\}$ is the body force and $\{P\}_0$ the surface 198 traction. Ω_{0_e} is a given element and $\Gamma_{0_e}^c$ is the potential contact surfaces, 199 both at the reference configuration. $[\mathbb{B}_L]$ is a matrix that depends on the 200 shape function derivatives and that converts the nodal displacements to the 201 linear part of the strains. Similarly, $[\mathbb{B}_{NL}]$ depends on the shape function 202 derivatives and relates the nodal displacements to the nonlinear portion of 203 the strains. $[\mathbb{D}]$ represents the material stiffness matrix and $[\mathbb{S}]$ stores the 204 second Piola-Kirchhoff components in diagonal blocks. \mathbf{n} and \mathbf{m} are normal 205 and tangential directions defined locally in each element. 206

207 2.7. Solution Technique

The solution at each time step is obtained via an incremental iterative scheme using (15). Starting from a known solution $(\{U\}^n)$, an increment

²¹⁰ $\Delta{U} = {U}^{n+1} - {U}^n$ is sought so that the residual forces $({R})$ are zero, ²¹¹ that is:

$$\{R\} = \{F\}_{ext}^{t+\Delta t} - \{F\}_{int}^{t} + \{F\}_{c}^{t+\Delta t} = 0.$$
 (21)

Since several sources of nonlinearity are present in the model, the residual rarely vanishes after the first iteration. This gives rise to unbalanced forces from the previous iteration (k) which are used to calculate nodal displacement corrections, so that:

$$([\mathbb{K}_L] + [\mathbb{K}_{NL}])^k \Delta \{U\}^{k+1} = \{R\}^k.$$
(22)

Iterations continue until the residual becomes negligible to within a specified tolerance. In this situation the converged nodal displacement $\{U\}^{n+1}$ is obtained by adding all the corrections:

$$\{U\}^{n+1} = \{U\} + \Delta\{U\}, \tag{23}$$

$$\Delta\{U\} = \sum_{k=1}^{NT} \Delta\{U\}^k, \tag{24}$$

where *NT* denotes the total number of iterations required for convergence. This process belongs to the class of Newton-based methods. In the present study, the computational model was implemented using Ansys Mechanical APDL 2021R2 solver.

223 3. Results and Discussion

224 3.1. Opening and Closing mechanism

Figure 9 shows the colour map representing the magnitude of the dis-225 placements observed both at opening and closing. The peculiar shape of the 226 leaflets makes the closing mechanism essentially composed of two global re-227 sponses: bending and torsion. Bending occurs in each of the leaflets so that 228 the regions which are farthest from the supports have greater movement. The 229 displacement amplitude shown by the free surfaces determines the degree of 230 opening of the valve (nearly 4 mm or half the nominal valve height). Torsion, 231 on the other hand, appears as a joint effect promoted by contact and which 232 leads to rotation of the upper circular boundaries near the valve axis. The 233 maximum displacements in the vertical downward direction is close to 30%234 of the nominal height. 235

The effects related to displacements are also observed in the stress field, as 236 shown in Figure 10. Although Von Mises stresses are used to predict yielding 237 in metallic materials, they also serve as a global analysis parameter because 238 they can represent a single scalar measure of the stress state. The colour map 230 in Figure 10 indicates the path of the internal stresses that maintain equilib-240 rium in both the open and closed position. The stresses are not distributed 241 regularly, but instead tend to follow the regions with greater vertical and 242 horizontal stiffness (near the supports and contacts). There is stress concen-243 tration in the joint region between the valve base and the vertical supports. 244 as expected, since it is a geometric singularity point. Points such as these 245 may require special attention in studies involving fatigue endurance. 246

247

Note that the imposed pressure load does not allow the full opening of



Figure 9: Global displacement field: (a) step = 0.5, (b) step = 1.0, (c) step = 2.5 and (d) step = 3.0. Units in mm.



Figure 10: Von Mises stress field: (a) step = 0.5, (b) step = 1.0, (c) step = 2.5 and (d) step = 3.0. Units in MPa.

the leaflets. The area of the circular base is $3.14 \ cm^2$. In the position of maximum opening, the minimum orifice area is $0.34 \ cm^2$, or approximately 11%. Naturally this opening rate can be changed by appropriate choice of materials and manufacturing thickness of the leaflets.

252 3.2. Distribution of bearing forces

An important aspect in any valve concept is to design supports capable of upholding the loads that will act on it, including the possibility of changing the modulus, direction and sense of the support reactions. This structural behaviour can be estimated based on the force fields distribution along the surfaces, in particular, near regions of discontinuity such as contacts and connections.

The way the forces are distributed along the annulus, as well as its proper geometric positioning can have an impact on the efficiency of the valve's operation. For instance, recent studies using in silico modelling have demonstrated that annuloplasticity (for natural mitral valves) with moderate annular reduction may be efficient for achieving optimal coaptation, compared with traditional annuloplasticity techniques [29].

To illustrate the internal force field, Figure 11 displays the fields for two 265 key step loads (1.0 and 3.0). These forces are derived from elemental data 266 (such as stresses and strains) calculated at integration points and then ex-267 trapolated to the nodes. The arrows indicate the direction of the acting 268 forces, and their sizes are in proportion. As can be seen, there is a concen-269 tration of forces on the contact surface which acts by mobilising the valve 270 opening and closing mechanisms. These forces are responsible for maintain-271 ing balance together with the forces on the support frame, represented by 272



²⁷³ the annular base and the vertical supports.

Figure 11: Internal vector force field: (a) step = 1.0 and (b) step = 3.0. Units in Newtons.

In Figure 12, the same force field is shown restricted to the support re-274 gions, adding transparency on the leaflet surfaces. This field of reaction 275 vectors confirms the existence of concentrated forces at the junction between 276 the valve base and the vertical supports (discussed in the Section 3.1). The 277 forces on both the circular base and the vertical supports induce bending. 278 For vertical ones the forces tend to bend the supports towards the interior of 279 the valve, while the base tends to bend in an alternating direction, either in-280 wards or outwards from the circle. As the forces lie in planes not necessarily 281 parallel to the shell surfaces, this indicates that the assumptions associated 282 with shell kinematics are preferable to those for membranes, at least for this 283 application. 284

These results illustrate that the supporting structure of the leaflets must be sufficiently resistant to bending forces in multiple directions for equilibrium to be preserved. In general the frames are made of rigid material, but



Figure 12: Internal force field along the supports: (a) step = 1.0 and (b) step = 3.0. Units in Newtons.

the possibility of more flexible structures is not excluded.

289 3.3. Effect of height for global stability

Three geometries with different heights (10 mm, 15 mm, and 20 mm) were analysed. The other model parameters such as ellipticity degree, constitutive model, contact assumptions, pressure loading curve, characteristic mesh size and kinematic boundary conditions were kept unchanged.

Table 2 shows a comparison between results for three designs. It can 294 be seen that the pattern of both leaflet opening and closing mechanisms 295 remains unchanged with increasing height/diameter ratio. The most notable 296 difference consists of the amplitude of displacements of the free surface at 297 the upper edge, which naturally tends to move further downwards. It is also 298 observed that the free hole of the opening (top view) increases with increasing 299 height of the valve. This gain in flexibility in the upper part can be explained 300 by the transfer of forces from the lower supports to the upright ones, leaving 301



Table 2: Maximum opening and closing positions considering three height (h) to diameter (d) ratios. In all cases the ellipticity degree is $\gamma = 1$. (Units in millimetres, reaction forces in Newtons, Stresses in MPa).

the surface less strained when opening. The alternating bending forces on the supports are maintained, but with different intensities, decreasing as the height increases.

305 3.4. Comparing S-shaped and U-shape design

Since the formulation proposed in this study is used to design S-shaped valves, it is interesting to compare the mechanical behaviour between the proposed leaflets and the traditional design (chosen here as the U-shaped). In this numerical experiment, both valves (Figure 13) have the same characteristic lengths, same boundary conditions and are simulated using the same constitutive material (Table 1). The contour plots representing mechanical indicators are showed in Table 3.



Figure 13: Perspective view of compared designs: (a) S-shape-based and (b) U-shapebased valves. In both cases, annular diameter = 20 mm, height = 10 mm, thickness of leaflets = 250 μ m are assumed.

In terms of global displacements, the S-shape design shows higher magnitudes compared to the U-shape. This effect is due to the larger portion

of material placed near the centre of the valve that can be moved vertically. This implies that for small pressure increments large displacements occur before tensile stresses can be mobilised to increase geometric rigidity. This arrangement contrasts with the U-shape. In mechanical terms, the U-shape is more stiff with respect to the forces that tend to close the valve.

Regarding the stress distribution, as expected, there is a concentration near the bottom of the vertical posts in the case of the S-shaped valve, which does not occur for the U-shape. This is due to the folding effect of the leaflets on themselves, which tends to create high local stresses. In open position, the S shape creates relatively comparative low stress regions, which provides the possibility for geometric optimisation in future designs.

These results show that the U-shaped valve is uniformly stressed in terms of bending, tensile and geometric stiffness. This doesn't leave much room for improvement on these surfaces. In the case of the S-shape, regions of high stress and regions of low stress can be seen, which may suggest that there exist more convenient curvatures to be used. This type of study could be carried out using parametric analyses and the equations presented here are very attractive for iterative search algorithms.

It is worth noting that the higher geometric stiffness of the S shape implies that the valve has a smaller central opening. This behaviour can induce flow direction as long as the leaflets can withstand the bending stresses. This characteristic can be exploited for specific purposes, in particular for irregular cavities to optimize the outflow volume. In this case, more in depth studies involving fluid-structure interaction are indicated.



Table 3: Results for S- and U- shape designs. For S-shape, the ellipticity degree is $\gamma = 1$. (Maximum opening and closing positions units in millimetres, Stresses in MPa).

339 4. Conclusion

In this work, based on the Wheatley concept, we presented a mathematical description of the leaflets of mitral valves and a nonlinear model of their movement. This procedure has been shown to be efficient for the mechanical performance study of a class of S-shaped leaflet designs. Some points are worth highlighting.

Regarding the purely geometric aspects, it can be remarked that:

- (i) This is a natural development from previous work on the Wheatley aortic valve ([19], [20] and [21]). We have taken this further by using ellipses to produce infinitely many designs controlled by a factor γ , which essentially measures the degree of ellipticity in the design;
- (ii) Two examples of the closed valve have been displayed for different 'ellipticities'. In addition the original Wheatley valve is shown both closed
 and partially open. Although these motions are kinematically admissible, they do not follow the equilibrium and material compatibility.
 However, these equations serve for both analysis and manufacturing
 design. Hence their importance.

³⁵⁶ Concerning the nonlinear structural model, it can be observed that:

- (iii) It was designed to predict the behaviour of a valve obtained from the equations assuming $\gamma = 1$. However the same procedure remains valid for different degrees of ellipticity;
- (iv) The essential aspects of this model are: linear behaviour between Green's
 deformations and the Second Piola-Kirchhoff Tensor, disregard of iner-

tial forces, consideration of large displacements, uniform pressure gradients on the surface and contact between leaflets. These aspects are considered to be the minimum necessary to capture the opening and closing motions. Changes in these assumptions necessarily imply an adaptation of the model presented here;

- (v) The S-shape of the leaflets induces internal forces that mobilise two
 global responses: bending and torsion. These two effects depends on
 the contact between the flexible surfaces and upon the stiffness of the
 support;
- (vi) The supports here have been considered completely rigid. In this situation, to maintain the valve equilibrium, the supports are subjected to
 forces in alternate directions. If flexible bearing frames were adopted
 these effects should be carefully considered in order to avoid structural
 instabilities.
- (vii) The mechanism for transmitting the equilibrium forces does not change
 when the height is increased (but not more than double) while maintaining the base diameter.
- With regard to the numerical experiment comparing an S-shaped valve with a U-shaped valve:
- (viii) The simulation showed that S-shaped leaflets have higher levels of bending, particularly in the central portion of the valve, when compared
 to U-shaped leaflets. However, this displacement does not hinder the
 closing mechanism, which remains stable and able to bear the required
 pressure.

In general, both the method used and the model proposed can be extended to study the Wheatley mitral valve in different scenarios.

388 Competing interests

The authors declare that David J. Wheatley is the inventor of the valve which was the subject of this study, holding patents in several countries including US Patent No. 9259313 2014, UK Patent No. EP2982340 2016, Europe Patent No. EP2979666 2016 and China Patent No. CN103384505. The full list is available at https://www.wheatleyresearch.co.uk/blog/developmenthistory.

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³⁹⁷ Ethical approval

³⁹⁸ Not required.

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Title of Paper: Mathematical representation and nonlinear modelling of the Wheatley mitral valve

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Conflicts of Interest

The authors declare that David J. Wheatley is the inventor of the valve which was the subject of this study, holding patents in several countries including US Patent No. 9259313 2014, UK Patent No. EP2982340 2016, Europe Patent No. EP2979666 2016 and China Patent No. CN103384505. The full list is available at https://www.wheatleyresearch.co.uk/blog/development-history

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