

ON THE ENERGY-MOMENTUM OF LIGHT: AN ALL OPTICAL VIEW OF THE ABRAHAM-MINKOWSKI CONTROVERSY Alexander J. Macleod* & Adam Noble

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Momentum of light in medium

The momentum of light in the linear vacuum is an unambiguous concept. Due to the simple constitutive relations $\mathbf{D} = \varepsilon_0 \mathbf{E}$, and $\mathbf{H} = \mathbf{B}/\mu_0$, the momentum density can be expressed as either $\mathbf{g}_{vac} = \mathbf{D} \times \mathbf{B}$ or $\mathbf{g}_{vac} = \mathbf{E} \times \mathbf{H}/c^2.$

Minkowski



Fig. 1: Atom of mass m and transition frequency ω_0 moving away from light source with velocity v. Beam frequency ω . System is in a medium of refractive index n. Atom will absorb photons from field if matching occurs between the Doppler-shifted and transition frequencies,

$\omega_0 \approx \omega \left(1 - \frac{nv}{c}\right).$

Abraham

PA



Fig. 2: A photon with energy $\hbar\omega$ travels in the z-direction with speed c towards a transparent block at rest (mass M, thickness L, refractive index n).

Photon speed becomes c/n in the medium, and uniform motion of centre of mass gives block displacement,

The presence of a medium with polarization \mathbf{P} and magnetization \mathbf{M} changes the constitutive relations,

$$\mathbf{D} = \varepsilon_0 \mathbf{E} + \mathbf{P}, \quad \mathbf{H} = \frac{1}{\mu_0} \mathbf{B} - \mathbf{M}.$$

There are now two possibilities for the momentum density, due to Minkowski and Abraham (see review [1]).

After absorption the atom will have a velocity v' and conservation of energy and momentum gives, $\frac{1}{2}mv'^2 + \hbar\omega_0 = \frac{1}{2}mv^2 + \hbar\omega, \qquad mv' = mv + p_{photon}.$ These equations can then be solved to give [2], $p_{photon} = \frac{\hbar\omega n}{c} \frac{2v}{v+v'} \approx \frac{\hbar\omega n}{c} \equiv p_M.$ This corresponds to momentum density, $\mathbf{g}_M = \mathbf{D} \times \mathbf{B}.$

 $\Delta z = (n-1)L\frac{\hbar\omega}{Mc^2}.$

Block must have acquired momentum from photon,

$$p_{block} = M \frac{\Delta z}{L(n/c)} = \left(1 - \frac{1}{n}\right) \frac{\hbar \omega}{c}.$$

Global conservation of momentum gives the total momentum of the system as $\hbar\omega/c$ and so [2],

$$p_{photon} = \frac{h\omega}{cn} \equiv p_A$$

This corresponds to momentum density, $\mathbf{g}_A = \mathbf{E} \times \mathbf{H}/c^2$.

A possible resolution?

Both forms of the momentum in a medium have experimental and theoretical support. It has been postulated that the total energy-momentum in either case must be the same, and that obtaining Abraham or Minkowski is dependent on how we partition the total energymomentum between the light and the medium. If the energy-momentum of the wave is $\Theta^{\mu}{}_{\nu}$ for Minkowski, and $\Omega^{\mu}{}_{\nu}$ for Abraham, then this can be expressed as,

 $\Theta^{\mu\nu} + T^{\mu\nu}_{M,\text{material}} \stackrel{?}{=} \Omega^{\mu\nu} + T^{\mu\nu}_{A,\text{material}}.$

Calculating the energy-momentum for real materials is extremely difficult to achieve. Can we approach this from a different perspective?

Non-linear vacuum electrodynamics

Spontaneously created and annihilated particle-antiparticle pairs in the quantum vacuum can become polarized by the presence of a strong electromagnetic field. This can cause the vacuum to exhibit non-linear properties analogous to that of a dielectric **material** such as acquiring an effective refractive index (see e.g. review [3]).

A probe field propagating on a strong background electromagnetic field gives us an



Fig. 3: Vacuum fluctuation.

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Energy-momentum tensor for non-linear electrodynamics

Non-linear vacuum electrodynamics is defined by the class of Lagrangian densities of the form $\mathcal{L}(X,Y)$ where $X = -F^{\mu\nu}F_{\mu\nu}/4$ and $Y = -F^{\mu\nu}F_{\mu\nu}/4$ are the only Lorentz invariant contractions of the electromagnetic field $F^{\mu\nu}$. This ensures that such theories maintain Lorentz invariance.

Using variational principles the full energy-momentum tensor is found to be,

$$T^{\mu}_{\ \nu} = H^{\mu\alpha}F_{\nu\alpha} + \delta^{\mu}_{\nu}\mathcal{L}.$$

The function $H^{\mu\nu}$ is the excitation tensor which describes the *constitutive relations*. In non-linear vacuum electrodynamics this is defined by,

 $H^{\mu\nu} \equiv \frac{\partial \mathcal{L}}{\partial X} F^{\mu\nu} + \frac{\partial \mathcal{L}}{\partial V} \widetilde{F}^{\mu\nu}.$

Equivalence of the full energy-momentum in Abraham and Minkowski

The energy-momentum tensor of the probe in Minkowski theory is given by, $\Theta^{\mu}{}_{\nu} = h^{\mu\alpha} f_{\nu\alpha} - \frac{1}{4} \delta^{\mu}_{\nu} h^{\alpha\beta} f_{\alpha\beta},$

To compare more directly with the light-matter case, we separate the electromagnetic field into a strong background $\mathcal{F}_{\mu\nu}$ and weak probe $f_{\mu\nu}$,

 $F_{\mu\nu} = \mathcal{F}_{\mu\nu} + f_{\mu\nu},$

such that the excitation tensor can be expanded as,

 $H^{\mu}{}_{\nu} = \mathcal{H}^{\mu}{}_{\nu} + h^{\mu}{}_{\nu} + \mathscr{H}^{\mu}{}_{\nu} + \mathcal{O}(f^3).$

Here $\mathcal{H}^{\mu}{}_{\nu} = O(f^0), h^{\mu}{}_{\nu} = O(f^1)$ and $\mathscr{H}^{\mu}{}_{\nu} = O(f^2)$. This can then be used to expand the full energy-momentum tensor and compare this with Abraham and Minkowski.

Abraham or Minkowski?

In the usual context of light-matter interactions, the velocity field u^{μ} in the Abraham formulation is the 4-velocity of the medium.

However, in non-linear vacuum electrodynamics, there is no physical medium present and no clear indication of how to interpret u^{μ} .

and in Abraham theory is,

 $\Omega^{\mu}{}_{\nu} = \Theta^{\mu}{}_{\nu} + \frac{1}{2} (f^{\mu\alpha}h_{\nu\alpha} - h^{\mu\alpha}f_{\nu\alpha}) - \frac{1}{2} \left[u^{\mu}u_{\alpha}(f_{\nu\beta}h^{\alpha\beta} - h_{\nu\beta}f^{\alpha\beta}) + u_{\nu}u^{\alpha}(f^{\mu\beta}h_{\alpha\beta} - h^{\mu\beta}f_{\alpha\beta}) \right]$ where u^{μ} ($u^2 = 1$) is a 4-velocity field.

A key result of this work is that the *exact* energy-momentum tensors of Abraham and Minkowski are equal. This then allows us to express the non-linear vacuum energymomentum as [4],

 $T^{\mu}_{\ \nu} = \Theta^{\mu}_{\ \nu} + \mathscr{M}^{\mu}_{\ \nu} = \Omega^{\mu}_{\ \nu} + \mathscr{A}^{\mu}_{\ \nu}.$

In non-linear vacuum electrodynamics we can obtain Abraham or Minkowski by appropriate partitioning of the total energy-momentum.



This means an arbitrary u^{μ} must be chosen, which introduces a preferred frame and breaks Lorentz invariance.

Minkowski energy-momentum is the correct description in non-linear vacuum electrodynamics.

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[3] Di Piazza, A., Müller, C., Hatsagortsyan, K., and Keitel, C. Rev. Mod. Phys. 84(3), 1177 (2012).

[4] Macleod, A. J., Noble, A., and Jaroszynski, D. A. Proc.SPIE **10234**, 0F (2017).