Mobility Analysis of Generalized Angulated Scissor-like Elements with the Reciprocal Screw Theory

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Abstract: This paper deals with the planar closed loop linkage consisting of a series of scissor-like elements connected by revolute joints. Because every generalized angulated element (GAE) subtends a constant angle during the motion, every angulated link was assumed as a PRRP linkage which has two prismatic joints and two revolute joints. Therefore, the two PRRP linkages of the GAE are individually movable with a single degree of freedom. The mobility of two types of GAEs was investigated with the method based on the screw theory. It has been proven that both types of GAEs are movable because the terminal constraints exerted to the common joint by the two linkages are equal.

Keywords: kinematics; scissor-like element; linkages; screw theory

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1. Introduction

Foldable structures, which can be transformed from a closed configuration to a expanded form, are developed from the earliest modern stages with Pinero's movable theater in 1961 [1]. Especially during the last decade, the retractable roof structures are widely used in long-span space structures [2]. On the other hand, with the further development of aerospace science and technology, deployable structures are of even greater interest in aerospace industries [3, 4].

Generally, the deployable structure based on the scissor-like element is the most commonly used form in engineering [5-8]. A scissor-like element in two dimensional forms is shown in Fig.1 (a). The axis of the revolute joint is perpendicular to the plane of the structure. The element has one degree-of-freedom and can be folded and deployed freely. Besides the foldable truss and dome structures [9], Zhao *et al.* [10, 11] proposed a foldable stair, which consists of a number of identical deployable scissor-like elements. In the early 1990's, Hoberman [12, 13] invented and patented a method for constructing loop assemblies formed by the modified scissor-like element, which consists of a pair of identical angulated rods connected by a revolute joint as shown in Fig.1 (b). Therefore, the unit is also called angulated scissor element or Hoberman's unit. You and Pellegrino [14] noted that the unit subtends a constant angle as their rods rotate while maintaining the end pivots on parallel lines. Thus the Hoberman's unit can create a closed-loop mechanism, which is called Hoberman's Linkage as shown in Fig.2. They also proposed two types of generalized angulated elements (GAEs), which subtend a constant angle during folding, but afford much

greater freedom of shapes than Hoberman's unit. The triangles of the two angulated rods of Type I GAE are isosceles triangles, i.e. $AE=DE$ and $BE=CE$, as shown in Fig.3. The triangles of the two angulated rods of Type II GAE are similar triangles, i.e. $AE/BE=DE/CE$ and $\angle AEB=\angle DEC$. As shown in Fig.4. the lengths of BE and AE are p and q . For Type II GAE, $AE/BE=DE/CE=w$. Then the lengths of CE and DE are wq and wq . The planar closed loop linkages which include Type I GAEs and Type II GAEs are given in Fig.5.

The focus of this paper is on the mobility analysis of this planar closed loop linkage. It is pertinent to point out that the kinematic treatment of linkages is receiving increasing attention in recent years [15-18]. Langbecker [5] proposed the foldability equation of deployable scissor structures, which is formulated using a purely geometric approach. Patel and Ananthasuresh [16] presented a kinematic theory behind Hoberman's and other inventions related to planar, radially foldable linkages. They showed that known types of foldable linkages can be derived by using a simple algebraic equation and even simpler design criteria. Based on the corresponding mobility conditions, Mao *et al.* [17] produced mobile double chains with both even and odd number of intersecting scissor-like pairs. A numerical algorithm was used by Nagaraj *et al.* [18] to evaluate the degree-of-freedom of pantograph masts by obtaining the null space of a constraint Jacobian matrix. Moreover, the principles of screw theory can also be used to study the degree-of-freedom of mechanism. Dai *et al.* [19-21] studied the mobility of a general type of foldable mechanisms using the screw theory. Zhao *et al.* [22] applied the screw theory to study the degree-of-freedom of simple

planar linkage and the mechanism theory of forming the spatial deployable units utilized in flat, cylindrical and spherical deployable structures. But they only focused on the scissor-like elements with straight rods. Then Cai *et al.* [23] extend this theory to the mobility analysis of the Hoberman's angulated unit. Dai *et al.* [24] and Wei *et al.* [25, 26] also used the reciprocal screws to study the mobility of the Hoberman structure and its variant. However, the scissor-like elements studied in the previous literatures are very simple and the geometry of these elements is symmetry. Two types of GAEs proposed by You and Pellegrino [14], which are not considered in these literatures, have asymmetry geometry and complex kinematic behavior. The mobility of GAEs will be studied in this paper with the screw theory. Prior to probing the mechanism theory of the GAE in Section 3, the kinematic description of GAEs will be given in Section 2. Section 4 draws conclusions.

2. Kinematic description of GAEs

As mentioned earlier, the main attribute of the planar closed loop linkages studied in this paper is that every angulated element subtends a constant angle during the motion. A simple generalized angulated element is shown in Fig.6a. The common revolute joint connects the two angulated rods of equal configuration. Therefore, points A, B, C and D can be seen as to be constrained to move along either of the dash lines P_1 and P_2 . Then we can interpret an angulated link as a PRRP linkage shown in Fig. 6(b), which has two prismatic joints and two revolute joints. Thus the two PRRP linkages of the GAE are individually movable with a single degree of freedom. Therefore, the GAE will be

movable only when the two linkages share the same coupler curve at point E [16].

3. Kinematic analysis of GAEs

The GAE can be considered as a parallel mechanism. It consists of the linkage ABE and linkage CDE, which are connected by a common joint E. Its mobility can be analyzed by the method based on the screw theory, as proposed in Ref. [24].

It is assumed that the subtended angle between line AD and line BC is denoted by 2α . The orthogonal coordinate oxy is shown in Fig. 7. The axis ox is the bisector of angle $\angle BOD$, and the axis oy is perpendicular to the ox axis in the mechanism plane. Now, the kinematics of the link ABE is studied firstly.

There are a revolute joint and a slide joint at point A. Then the twists of the RP chain are given as

$$\Phi_A = [\Phi_{A_R} \quad \Phi_{A_S}] \quad (1)$$

where

$$\begin{aligned} \Phi_{A_R} &= \{0 \quad 0 \quad 1 \quad y_A \quad -x_A \quad 0\}^T \\ \Phi_{A_S} &= \{0 \quad 0 \quad 0 \quad \cos \alpha \quad -\sin \alpha \quad 0\}^T \end{aligned}$$

and x_A, y_A are the coordinates of point A.

According to the screw theory, the terminal constraint screw matrices of the RP chain are immediately obtained by solving the reciprocal screw equations

$$\Phi_A^T \mathbf{E} \Gamma_A = 0 \quad (2)$$

where $\mathbf{E}$ is the swap operator and its formulation is expressed as follows

$$\mathbf{E} = \begin{bmatrix} 0 & \mathbf{I}_3 \\ \mathbf{I}_3 & 0 \end{bmatrix} \quad \text{and} \quad \mathbf{I}_3 = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}.$$

Therefore, the terminal constraints external to link ABE by the RP chain are

$$\mathbf{\Gamma}_A = \begin{bmatrix} \mathbf{\Gamma}_{A,1} & \mathbf{\Gamma}_{A,2} & \mathbf{\Gamma}_{A,3} & \mathbf{\Gamma}_{A,4} \end{bmatrix} \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ f_4 \end{bmatrix} \quad (3)$$

where f_i ($i=1, 2, 3, 4$) represent the tensities of the screws and

$$\begin{aligned} \mathbf{\Gamma}_{A,1} &= \{-\sin \alpha \quad -\cos \alpha \quad 0 \quad 0 \quad 0 \quad -x_A \cos \alpha + y_A \sin \alpha\}^T \\ \mathbf{\Gamma}_{A,2} &= \{0 \quad 0 \quad 1 \quad 0 \quad 0 \quad 0\}^T \\ \mathbf{\Gamma}_{A,3} &= \{0 \quad 0 \quad 0 \quad 1 \quad 0 \quad 0\}^T \\ \mathbf{\Gamma}_{A,4} &= \{0 \quad 0 \quad 0 \quad 0 \quad 1 \quad 0\}^T \end{aligned}.$$

Similarly, the terminal constraints external to link ABE by the RP chain at point B are

$$\mathbf{\Gamma}_B = \begin{bmatrix} \mathbf{\Gamma}_{B,1} & \mathbf{\Gamma}_{B,2} & \mathbf{\Gamma}_{B,3} & \mathbf{\Gamma}_{B,4} \end{bmatrix} \begin{bmatrix} f_5 \\ f_6 \\ f_7 \\ f_8 \end{bmatrix} \quad (4)$$

where

$$\begin{aligned} \mathbf{\Gamma}_{B,1} &= \{\sin \alpha \quad -\cos \alpha \quad 0 \quad 0 \quad 0 \quad -x_B \cos \alpha - y_B \sin \alpha\}^T \\ \mathbf{\Gamma}_{B,2} &= \{0 \quad 0 \quad 1 \quad 0 \quad 0 \quad 0\}^T \\ \mathbf{\Gamma}_{B,3} &= \{0 \quad 0 \quad 0 \quad 1 \quad 0 \quad 0\}^T \\ \mathbf{\Gamma}_{B,4} &= \{0 \quad 0 \quad 0 \quad 0 \quad 1 \quad 0\}^T \end{aligned},$$

f_i ($i=5, 6, 7, 8$) represent the tensities of the screws and x_B, y_B are the coordinates of point B.

We can find the terminal constraints exerted to link ABE using Eqs. (3) and (4) as

$$\mathbf{\Gamma}_{ABE} = \begin{bmatrix} \mathbf{\Gamma}_{A,1} & \mathbf{\Gamma}_{B,1} & \mathbf{\Gamma}_{A,2} & \mathbf{\Gamma}_{A,3} & \mathbf{\Gamma}_{A,4} \end{bmatrix} \begin{bmatrix} f_1 \\ f_5 \\ f_2 + f_6 \\ f_3 + f_7 \\ f_4 + f_8 \end{bmatrix} \quad (5)$$

Consequently, the twist of the link ABE can be solved with the reciprocal screw theory

$$\mathbf{\Phi}_{ABE}^T \mathbf{E} \mathbf{\Gamma}_{ABE} = 0 \quad (6)$$

Then the twist of the link ABE, denoted by $\mathbf{\Phi}_{ABE}$, is

$$\mathbf{\Phi}_{ABE} = \left\{ \begin{bmatrix} 0 \\ 0 \\ 1 \\ \frac{1}{2}[y_A + y_B + (x_B - x_A) \cot \alpha] \\ \frac{1}{2}[(y_A - y_B) \tan \alpha - (x_A + x_B)] \\ 0 \end{bmatrix} \right\} \quad (7)$$

The twist of the revolute joint at point E can be given as

$$\mathbf{\Phi}_{E_R} = \{0 \quad 0 \quad 1 \quad y_E \quad -x_E \quad 0\}^T \quad (8)$$

where x_E is the coordinate of point E in the ox direction.

Then the twist of joint E can be given as

$$\mathbf{\Phi}_E = [\mathbf{\Phi}_{E_R} \quad \mathbf{\Phi}_{ABE}] \quad (9)$$

The terminal constraints exerted to joint E by the link ABE can be obtained by solving the reciprocal screw equation

$$\mathbf{\Phi}_E^T \mathbf{E} \mathbf{\Gamma}_E = 0 \quad (10)$$

The terminal constraints are given as

$$\mathbf{\Gamma}_E = \begin{Bmatrix} \beta_1 & 0 & 0 & 0 \\ \beta_2 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \beta_3 & 0 & 0 & 0 \end{Bmatrix} \begin{bmatrix} f_9 \\ f_{10} \\ f_{11} \\ f_{12} \end{bmatrix} \quad (11)$$

where

$$\begin{cases} \beta_1 = x_E + \frac{1}{2}[(y_A - y_B) \tan \alpha - (x_A + x_B)] \\ \beta_2 = y_E - \frac{1}{2}[y_A + y_B - (x_A - x_B) \cot \alpha] \\ \beta_3 = -\frac{1}{2}\{y_E[(y_A - y_B) \tan \alpha - (x_A + x_B)] + x_E[y_A + y_B - (x_A - x_B) \cot \alpha]\} \end{cases} \quad (12)$$

Similarly, the terminal constraints exerted to joint E by the link CDE can be obtained

as

$$\mathbf{\Gamma}_{E'} = \begin{Bmatrix} \beta_1' & 0 & 0 & 0 \\ \beta_2' & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ \beta_3' & 0 & 0 & 0 \end{Bmatrix} \begin{bmatrix} f_{13} \\ f_{14} \\ f_{15} \\ f_{16} \end{bmatrix} \quad (13)$$

where

$$\begin{cases} \beta_1' = x_E + \frac{1}{2}[(y_D - y_C) \tan \alpha - (x_D + x_C)] \\ \beta_2' = y_E - \frac{1}{2}[y_D + y_C - (x_D - x_C) \cot \alpha] \\ \beta_3' = -\frac{1}{2}\{y_E[(y_D - y_C) \tan \alpha - (x_D + x_C)] + x_E[y_D + y_C - (x_D - x_C) \cot \alpha]\} \end{cases} \quad (14)$$

It can be found that the GAE will be movable if the terminal constraints exerted to joint E by the link ABE and the link CDE are equal, which is

$$\mathbf{\Gamma}_E = \mathbf{\Gamma}_{E'} \quad (15)$$

It can be seen from Fig.7 that the relations between the coordinates of the joints and the subtended angle α are

$$\begin{cases} \tan \alpha = \frac{y_B}{x_B} = \frac{y_C}{x_C} \\ \tan \alpha = -\frac{y_A}{x_A} = -\frac{y_D}{x_D} \end{cases} \quad (16)$$

3.1 Type I GAE

For Type I GAE, the triangle $\triangle ADE$ and $\triangle BCE$ are isosceles. For triangle $\triangle ADE$, F is the midpoint of AD, and then line EF is perpendicular to line AD. Thus its line equation is

$$y = \cot \alpha x + b \quad (17)$$

Substituting the coordinates of joint F in Eq. (17) leads to

$$y = \cot \alpha x - \frac{x_A + x_D}{\sin 2\alpha} \quad (18)$$

For triangle $\triangle BCE$, G is the midpoint of BC, and then line EG is perpendicular to line BC. Similarly, the equation of line EG can be given as

$$y = -\cot \alpha x + \frac{x_B + x_C}{\sin 2\alpha} \quad (19)$$

Thus the coordinates of joint E can be obtained by Eqs. (18) and (19) as

$$\begin{cases} x_E = \frac{x_A + x_B + x_C + x_D}{4} \frac{1}{\cos^2 \alpha} \\ y_E = \frac{x_B + x_C - x_A - x_D}{4} \frac{1}{\sin \alpha \cos \alpha} \end{cases} \quad (20)$$

Substituting Eqs. (16) and (20) into Eq. (12) leads to

$$\begin{cases} \beta_1 = \frac{x_C + x_D - x_A - x_B}{4} \sec^2 \alpha \\ \beta_2 = \frac{x_C - x_D + x_A - x_B}{2 \sin 2\alpha} \\ \beta_3 = \frac{x_A x_C - x_B x_D}{2 \sin 2\alpha \cos^2 \alpha} \end{cases} \quad (21)$$

The coefficients of Eq. (13) can be obtained by substituting Eqs. (16) and (20) into Eq. (14)

$$\begin{cases} \beta_1' = \frac{x_A + x_B - x_C - x_D}{4} \sec^2 \alpha \\ \beta_2' = \frac{x_B + x_D - x_A - x_C}{2 \sin 2\alpha} \\ \beta_3' = \frac{-x_A x_C + x_B x_D}{2 \sin 2\alpha \cos^2 \alpha} \end{cases} \quad (22)$$

Substituting Eqs. (21) and (22) into Eqs. (11) and (13) respectively, it can be found that the terminal constraints exerted to joint E by the link ABE and the link CDE are equal. That is to say the Type I GAE is movable.

3.2 Type II GAE

For Type II GAE, the geometric conditions are

$$\begin{cases} \frac{AE}{DE} = \frac{BE}{CE} \\ \angle AED = \angle BEC \end{cases} \quad (23)$$

Thus $\triangle ADE$ and $\triangle BCE$ are similar triangles which leads to

$$\frac{AE}{BE} = \frac{DE}{CE} = \frac{AD}{BC} \quad (24)$$

This can also be rewritten as

$$\frac{(y_E - y_A)^2 + (x_E - x_A)^2}{(y_E - y_B)^2 + (x_E - x_B)^2} = \frac{(y_E - y_D)^2 + (x_E - x_D)^2}{(y_E - y_C)^2 + (x_E - x_C)^2} = \frac{(y_A - y_D)^2 + (x_A - x_D)^2}{(y_B - y_C)^2 + (x_B - x_C)^2} \quad (25)$$

Then the coordinates of joint E can be obtained with Eqs. (16) and (25) as

$$\begin{cases} x_E = \frac{x_A x_C - x_B x_D}{x_A - x_B + x_C - x_D} \\ y_E = \frac{x_A x_C - x_B x_D}{x_A + x_B - x_C - x_D} \tan \alpha \end{cases} \quad (26)$$

Substituting Eqs. (16) and (26) into Eq. (12) leads to

$$\begin{cases} \beta_1 = \frac{1}{2} \frac{(2x_A x_C \cos^2 \alpha - 2x_B x_D \cos^2 \alpha + x_A x_D - x_A x_C - x_A^2 + x_B x_D + x_B^2 - x_B x_C)}{(x_A + x_C - x_B - x_D) \cos^2 \alpha} \\ \beta_2 = -\frac{1}{2} \frac{(2x_A x_C \cos^2 \alpha - 2x_B x_D \cos^2 \alpha + x_A x_D - x_A x_C - x_A^2 + x_B x_D + x_B^2 - x_B x_C)}{(x_A + x_B - x_C - x_D) \sin \alpha \cos \alpha} \\ \beta_3 = -\frac{1}{2} \frac{(x_A x_C - x_B x_D)(2x_A x_C \cos^2 \alpha - 2x_B x_D \cos^2 \alpha + x_A x_D - x_A x_C - x_A^2 + x_B x_D + x_B^2 - x_B x_C)}{(x_A + x_B - x_C - x_D)(x_A + x_C - x_B - x_D) \sin \alpha \cos \alpha} \end{cases}$$

(27)

The coefficients of Eq. (13) can also be obtained by substituting Eqs. (16) and (26) into Eq. (14)

$$\left\{ \begin{array}{l} \beta_1' = \frac{1}{2} \frac{(2x_A x_C \cos^2 \alpha - 2x_B x_D \cos^2 \alpha + x_B x_C - x_C^2 - x_A x_C + x_D^2 + x_B x_D - x_A x_D)}{(x_A + x_C - x_B - x_D) \cos^2 \alpha} \\ \beta_2' = -\frac{1}{2} \frac{(2x_A x_C \cos^2 \alpha - 2x_B x_D \cos^2 \alpha + x_B x_C - x_C^2 - x_A x_C + x_D^2 + x_B x_D - x_A x_D)}{(x_A + x_B - x_C - x_D) \sin \alpha \cos \alpha} \\ \beta_3' = -\frac{1}{2} \frac{(x_A x_C - x_B x_D)(2x_A x_C \cos^2 \alpha - 2x_B x_D \cos^2 \alpha + x_B x_C - x_C^2 - x_A x_C + x_D^2 + x_B x_D - x_A x_D)}{(x_A + x_B - x_C - x_D)(x_A + x_C - x_B - x_D) \sin \alpha \cos \alpha} \end{array} \right. \quad (28)$$

Combining and rearranging Eqs. (27) and (28) leads to

$$\left\{ \begin{array}{l} \frac{\beta_1}{\beta_2} = -\frac{x_A + x_B - x_C - x_D}{x_A + x_C - x_B - x_D} \tan \alpha \\ \frac{\beta_1}{\beta_3} = -\frac{x_A + x_B - x_C - x_D}{x_A x_C - x_B x_D} \sin \alpha \cos \alpha \end{array} \right. \quad (29)$$

$$\left\{ \begin{array}{l} \frac{\beta_1'}{\beta_2'} = -\frac{x_A + x_B - x_C - x_D}{x_A + x_C - x_B - x_D} \tan \alpha \\ \frac{\beta_1'}{\beta_3'} = -\frac{x_A + x_B - x_C - x_D}{x_A x_C - x_B x_D} \sin \alpha \cos \alpha \end{array} \right. \quad (30)$$

From Eqs. (29) and (30) it can be shown that

$$\beta_1 : \beta_2 : \beta_3 = \beta_1' : \beta_2' : \beta_3' \quad (31)$$

Substituting Eq. (31) into Eqs. (11) and (13), it can be found that the terminal constraints exerted to joint E by the link ABE and the link CDE are equal. That also concludes Type II GAE is movable.

4. Conclusions

In this paper, we studied the kinematics of generalized angulated scissor-like elements,

which forms the planar closed loop linkage. Firstly, every angulated rod was assumed as a PRRP linkage, which has two prismatic joints and two revolute joints. Thus the two PRRP linkages of the GAE are individually movable with a single degree of freedom. The mobility was analyzed with the method based on the screw theory. It has been shown that both types of GAE are movable because the terminal constraints imposed onto the common joint E by the link ABE and CDE are equal. It should be noted that a more general angulated scissor element containing intermediate parallelograms, as reported by You and Pellegrino [14], is yet to be proven with the help of this theory. Further extensions of this approach include the kinematic analysis of planar closed loop double chain linkages which satisfy the loop parallelogram condition [17] or other multi-loop linkages [27].

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(a)

(b)

Fig.1 Scissor-like elements.

(a)

(b)

Fig.2 Hoberman's Linkage.

Fig.3 Type I GAE formed by angulated rods

Fig.4 Type II GAE formed by angulated rods

(a) Type I

(b) Type II

Fig. 5 Planar closed-loop linkages with GAEs

(a)

(b)

Fig.6 Kinematic description of GAEs.

Fig.7 Geometric description of GAEs.